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Transverse-momentum and pseudorapidity distributions of charged hadrons in pp collisions at $\sqrt{s}=7$ TeV

CMS Collaboration ; Khachatryan, V ; Amsler, C ; De Visscher, S ; Chiochia, V ; et al

Abstract: Charged-hadron transverse-momentum and pseudorapidity distributions in proton-proton collisions at square root of $s = 7$ TeV are measured with the inner tracking system of the CMS detector at the LHC. The charged-hadron yield is obtained by counting the number of reconstructed hits, hit pairs, and fully reconstructed charged-particle tracks. The combination of the three methods gives a charged-particle multiplicity per unit of pseudorapidity $dN(\text{ch})/d|\eta| < 0.5 = 5.78 \pm 0.01(\text{stat}) \pm 0.23(\text{syst})$ for non-single-diffractive events, higher than predicted by commonly used models. The relative increase in charged-particle multiplicity from square root of $s = 0.9$ to 7 TeV is $[66.1 \pm 1.0(\text{stat}) \pm 4.2(\text{syst})]\%$. The mean transverse momentum is measured to be $0.545 \pm 0.005(\text{stat}) \pm 0.015(\text{syst})$ GeV/c. The results are compared with similar measurements at lower energies.

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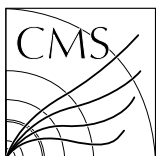
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Transverse-momentum and pseudorapidity distributions of charged hadrons in pp collisions at $\sqrt{s} = 7$ TeV

The CMS Collaboration*

Abstract

Charged-hadron transverse-momentum and pseudorapidity distributions in proton-proton collisions at $\sqrt{s} = 7$ TeV are measured with the inner tracking system of the CMS detector at the LHC. The charged-hadron yield is obtained by counting the number of reconstructed hits, hit-pairs, and fully reconstructed charged-particle tracks. The combination of the three methods gives a charged-particle multiplicity per unit of pseudorapidity $dN_{\text{ch}}/d\eta|_{|\eta|<0.5} = 5.78 \pm 0.01$ (stat.) ± 0.23 (syst.) for non-single-diffractive events, higher than predicted by commonly used models. The relative increase in charged-particle multiplicity from $\sqrt{s} = 0.9$ to 7 TeV is $(66.1 \pm 1.0$ (stat.) ± 4.2 (syst.))%. The mean transverse momentum is measured to be 0.545 ± 0.005 (stat.) ± 0.015 (syst.) GeV/ c . The results are compared with similar measurements at lower energies.

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*See Appendix A for the list of collaboration members

Introduction. Measurements of particle yields and kinematic distributions are an essential first step in exploring a new energy regime of particle collisions. Such studies contribute to our understanding of the physics of hadron production, including the relative roles of soft and hard scattering contributions, and help construct a solid foundation for other investigations. In the complicated environment of LHC pp collisions [1], firm knowledge of the rates and distributions of inclusive particle production is needed to distinguish rare signal events from the much larger backgrounds of soft hadronic interactions. They will also serve as points of reference for the measurement of nuclear-medium effects in Pb-Pb collisions in the LHC heavy ion program.

The bulk of the particles produced in pp collisions arise from soft interactions, which are modeled only phenomenologically. Experimental results provide the critical guidance for tuning these widely-used models and event generators. Soft collisions are commonly classified as elastic scattering, inelastic single-diffractive (SD) dissociation, double-diffractive (DD) dissociation, and inelastic non-diffractive (ND) scattering [2]. (Double-Pomeron exchange is treated as DD in this paper.) All results presented here refer to inelastic non-single-diffractive (NSD) interactions, and are based on an event selection that retains a large fraction of the ND and DD events, while disfavoring SD events.

The measurements focus on transverse-momentum p_T and pseudorapidity η distributions. The pseudorapidity, commonly used to characterize the direction of particle emission, is defined as $\eta = -\ln \tan(\theta/2)$, where θ is the polar angle of the direction of the particle with respect to the anti-clockwise beam direction. The count of primary charged hadrons, N_{ch} , is defined to include decay products of particles with proper lifetimes less than 1 cm. Products of secondary interactions are excluded, and a per-cent-level correction is applied for prompt leptons. The measurements reported here are of $dN_{\text{ch}}/d\eta$ and dN_{ch}/dp_T in the pseudorapidity range $|\eta| < 2.4$ and closely follow our previous analysis of minimum-bias data at lower centre-of-mass energies of $\sqrt{s} = 0.9$ and 2.36 TeV as reported in Ref. [3].

The data for this study are drawn from an integrated luminosity of $1.1 \mu\text{b}^{-1}$ recorded with the Compact Muon Solenoid (CMS) experiment [4] on March 30th, 2010, during the first hour of the LHC operation at $\sqrt{s} = 7$ TeV. These results are the highest centre-of-mass energy measurements of the $dN_{\text{ch}}/d\eta$ and dN_{ch}/dp_T distributions conducted at a particle collider so far and complement the other recent measurements of the ALICE experiment at 7 TeV [5].

Experimental methods. A detailed description of the CMS experiment can be found in Ref. [4]. The detectors used for the present analysis are the pixel and silicon-strip tracker (SST), covering the region $|\eta| < 2.5$ and immersed in a 3.8 T axial magnetic field. The pixel tracker consists of three barrel layers and two end-cap disks at each barrel end. The forward calorimeter (HF), which covers the region $2.9 < |\eta| < 5.2$, was also used for event selection. The detailed Monte Carlo simulation (MC) of the CMS detector response is based on Geant4 [6].

The event selection and analysis methods in this paper are identical to those used in Ref. [3], where more details can be found. The inelastic pp collision rate was about 50 Hz. At these rates, the fraction of events in the data, where two or more minimum-bias collisions occurred in the same bunch crossing, is estimated to be less than 0.3% and was neglected. Any hit in the beam scintillator counters (BSC, $3.23 < |\eta| < 4.65$) coinciding with colliding proton bunches was used for triggering the data acquisition. A sample mostly populated with NSD events was selected by requiring a primary vertex (PV) to be reconstructed with the tracker, together with at least one HF tower in each end with more than 3 GeV total energy. Beam-halo and other beam-background events were rejected as described in Ref. [3]. The remaining fraction of background events in the data was found to be less than 2×10^{-5} . The numbers of events

Table 1: Numbers of events passing the selection cuts. The selection criteria are applied in sequence, i.e., each line includes the selection from the previous ones.

Selection	Number of events
Colliding bunches + one BSC signal	68 512
Reconstructed PV	61 551
HF coincidence	55 113
Beam-halo rejection	55 104
Other beam-background rejection	55 100

Table 2: Fractions of SD, DD, ND, and NSD processes obtained from the PYTHIA and PHOJET event generators before any selection, and the corresponding selection efficiencies determined from the MC simulation.

	PYTHIA		PHOJET	
	Frac.	Sel. eff.	Frac.	Sel. eff.
SD	19.2%	26.7%	13.8%	30.7%
DD	12.9%	33.6%	6.6%	48.3%
ND	67.9%	96.4%	79.6%	97.1%
NSD	80.8%	86.3%	86.2%	93.4%

satisfying the selection criteria are listed in Table 1.

The event selection efficiency was estimated with simulated events using the PYTHIA [7, 8] and PHOJET [9, 10] event generators. The relative event fractions of SD, DD, and ND processes and their respective event selection efficiencies are listed in Table 2. The fraction of diffractive events is predicted by the models to decrease as a function of collision energy, while the selection efficiency increases. At $\sqrt{s} = 7$ TeV, the fraction of SD (DD) events in the selected data sample, estimated with PYTHIA and PHOJET, are 6.8% (5.8%) and 5.0% (3.8%), respectively, somewhat higher than at $\sqrt{s} = 0.9$ and 2.36 TeV [3]. With PYTHIA, the overall correction for the selection efficiency of NSD processes and for the fraction of SD events remaining in the data sample lowers the measured charged-particle multiplicity by 6% compared with the uncorrected distribution.

The $dN_{\text{ch}}/d\eta$ distributions were obtained, as in Ref. [3], with three methods, based on counting the following quantities: (i) reconstructed clusters in the barrel part of the pixel detector; (ii) pixel tracklets composed of pairs of clusters in different pixel barrel layers; and (iii) tracks reconstructed in the full tracker volume. The third method also allows a measurement of the dN_{ch}/dp_T distribution. All three methods rely on the reconstruction of a PV [11]. The PV reconstruction efficiency was found to be 98.3% (98.0%) in data (MC), evaluated after all other event selection cuts. In case of multiple PV candidates, the vertex with the largest track multiplicity was chosen. The three methods are sensitive to the measurement of particles down to p_T values of about 30, 50, and 100 MeV/c, respectively. Only 0.5, 1.5, and 5% of all charged particles are estimated to be produced below these p_T values, respectively, and these fractions were corrected for.

The measurements were corrected for the geometrical acceptance ($\approx 2\%$), efficiency ($\approx 5\text{--}10\%$), fake ($< 1\%$) and duplicate tracks ($< 0.5\%$), low- p_T particles curling in the axial magnetic field

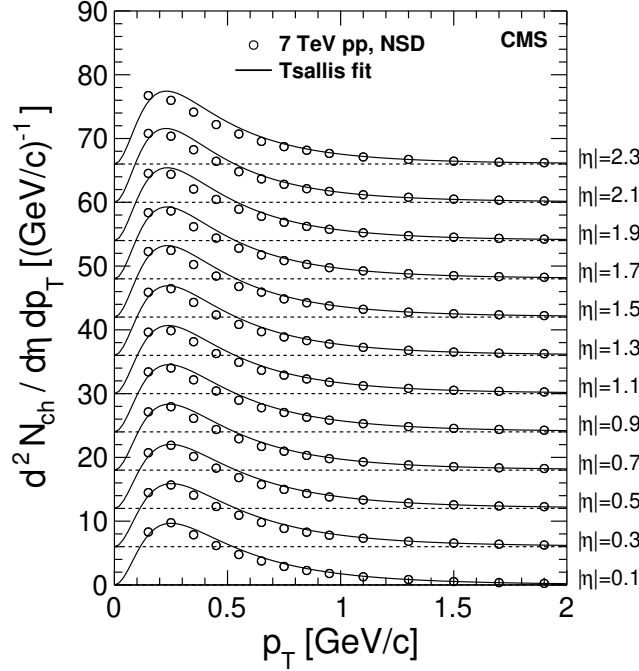


Figure 1: Differential yield of charged hadrons in the range $|\eta| < 2.4$ in 0.2-unit-wide bins of $|\eta|$ in NSD events. The solid curves represent fits of Eq. 1 to the data. The measurements with increasing η are successively shifted by six units along the vertical axis.

(<1%), decay products of long-lived hadrons (<2%) and photon conversions (<1%) and inelastic hadronic interactions in the detector material (≈ 1 -2%), where the size of the corrections in parentheses refer to the tracking method. The PYTHIA parameter set from Ref. [8] was chosen to determine the corrections, because it reproduces the $dN_{\text{ch}}/d\eta$ and charged-particle multiplicity distributions, as well as other control distributions at 7 TeV, better than other available tuning parameter sets. Although the corrections do not depend significantly on the model used, it is indeed important that the simulated data set contains a sufficient number of high-multiplicity events to determine these corrections with the desired accuracy.

Results. For the measurement of the dN_{ch}/dp_T distribution, charged-particle tracks with p_T in excess of 0.1 GeV/c were used in 12 different $|\eta|$ bins, from 0 to 2.4. The average charged-hadron yields in NSD events are shown in Fig. 1 as a function of p_T and $|\eta|$. The Tsallis parametrization [12–14],

$$E \frac{d^3 N_{\text{ch}}}{dp^3} = \frac{1}{2\pi p_T} \frac{E}{p} \frac{d^2 N_{\text{ch}}}{d\eta dp_T} = C \frac{dN_{\text{ch}}}{dy} \left(1 + \frac{E_T}{nT} \right)^{-n}, \quad (1)$$

where $y = 0.5 \ln[(E + p_z)/(E - p_z)]$, $E_T = \sqrt{m^2 + p_T^2} - m$, and m is the charged pion mass, was fitted to the data. The p_T spectrum of charged hadrons, $1/(2\pi p_T) d^2 N_{\text{ch}}/d\eta dp_T$, measured in the range $|\eta| < 2.4$, is shown in Fig. 2 for data at 0.9, 2.36, and 7 TeV. The high- p_T reach of the data is limited by the increase of systematic uncertainties with p_T . The fit to the data (Eq. 1) is mainly used for extrapolations to $p_T = 0$, but is not expected to give a good description of the data in all η bins with only two parameters. The parameter T and the exponent n were found to be $T = 0.145 \pm 0.005$ (syst.) GeV and $n = 6.6 \pm 0.2$ (syst.). The average p_T , calculated from a combination of the measured data points and the low- and high- p_T contributions as

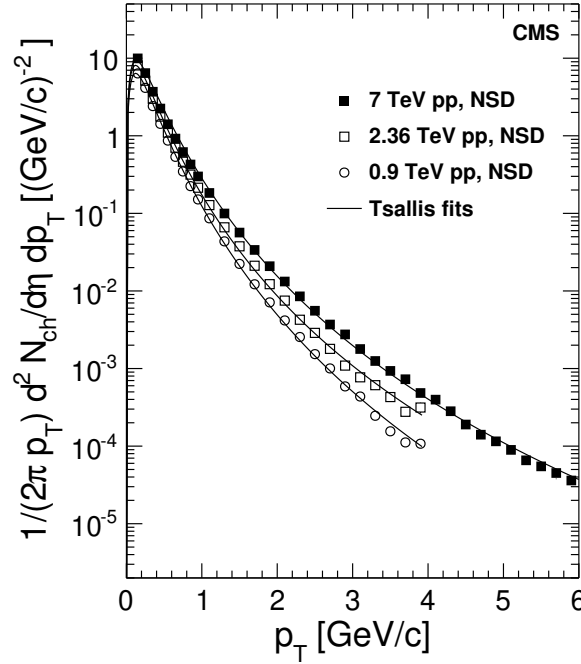


Figure 2: Charged-hadron yield in the range $|\eta| < 2.4$ in NSD events as a function of p_T ; the systematic uncertainties are smaller than the symbols. The measurements at $\sqrt{s} = 0.9$ and 2.36 TeV [3] are also shown. The solid lines represent fits of Eq. 1 to the data.

determined from the fit, is $\langle p_T \rangle = 0.545 \pm 0.005$ (stat.) ± 0.015 (syst.) GeV/c.

Experimental uncertainties related to the trigger and event selection are common to all the analysis methods. The uncertainty related to the presence of SD (DD) events in the final sample was estimated to be 1.4% (1.1%), based on consistency checks between data and simulation for diffractive event candidates. The total event selection uncertainty, which also includes the selection efficiency of the BSC and HF, was found to be 3.5%. Based on studies similar to those presented in Ref. [3], additional 3% and 2% uncertainties were assigned to the tracklet and track reconstruction algorithm efficiencies, respectively. Corrections at the percent level were applied to the final results to extrapolate to $p_T = 0$. The uncertainty on these extrapolation corrections was found to be less than 1%. All other uncertainties are identical to those listed in Ref. [3]. The $dN_{ch}/d\eta$ measurements were repeated on a separate data sample without any magnetic field, for which almost no p_T extrapolation is needed, and gave results consistent within 1.5%. The final systematic uncertainties for the pixel counting, tracklet, and track methods were found to be 5.7%, 4.6%, and 4.3%, respectively, and are strongly correlated.

For the $dN_{ch}/d\eta$ measurements, the results for the three individual layers within the cluster-counting method were found to be consistent within 1.2% and were combined. The three layer-pairs in the pixel-tracklet method provided results that agreed within 0.6% and were also combined. Finally, the results from the three different measurement methods, which agree with the combined result within 1% to 4% depending on η , were averaged. The final $dN_{ch}/d\eta$ distributions are shown in Fig. 3 for $\sqrt{s} = 0.9, 2.36$, and 7 TeV. The CMS results are compared with measurements made by other experiments. In the ATLAS Collaboration analysis [17], events and particles were selected in a different region of phase space, which makes a direct comparison difficult. Their results are therefore not included in the figure.

The results can also be compared to earlier experiments as a function of \sqrt{s} . The energy depen-

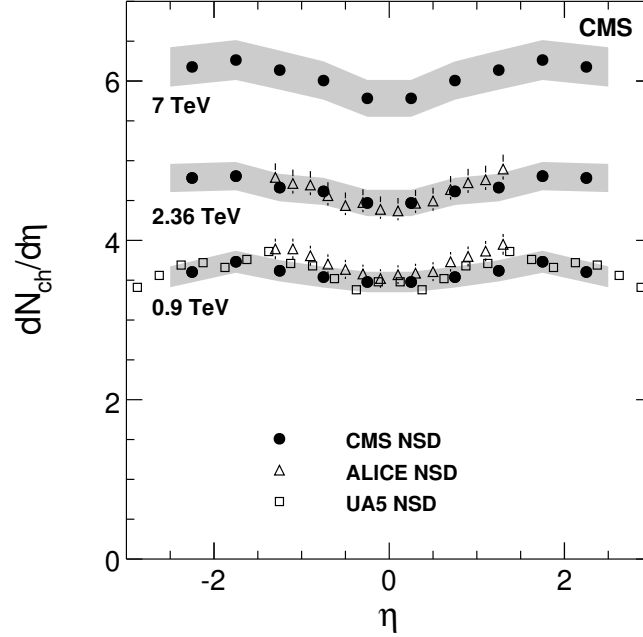


Figure 3: Distributions of $dN_{ch}/d\eta$, averaged over the three measurement methods and compared with data from UA5 [15] ($p\bar{p}$, with statistical errors only) and ALICE [16] (with systematic uncertainties). The shaded band shows systematic uncertainties of the CMS data. The CMS and UA5 data are averaged over negative and positive values of η .

dence of the average charged-hadron p_T can be described by a quadratic function of $\ln s$ [18]. As shown in Fig. 4, the present measurement follows this trend. The choice of the $|\eta|$ interval can influence the average p_T value by a few per cent.

For $|\eta| < 0.5$, the average charged multiplicity density is $dN_{ch}/d\eta = 5.78 \pm 0.01$ (stat.) ± 0.23 (syst.) for NSD events. The \sqrt{s} dependence of the measured $dN_{ch}/d\eta|_{\eta \approx 0}$ is shown in Fig. 5, which includes data from various other experiments. The $dN_{ch}/d\eta$ results reported here show a rather steep increase between 0.9 and 7 TeV, which is measured to be $(66.1 \pm 1.0$ (stat.) ± 4.2 (syst.))%. Using a somewhat different event selection, the ALICE collaboration has found a similar increase of $(57.6 \pm 0.4$ (stat.) $^{+3.6}_{-1.8}$ (syst.))% [5]. The measured charged-particle multiplicity is accurate enough to distinguish among most sets of event-generator tuning parameter values and various models. The measured value at 7 TeV significantly exceeds the prediction of 4.57 from PHOJET [9, 10], and the predictions of 3.99, 4.18, and 4.34 from the DW [22], ProQ20 [23], and Perugia0 [24] tuning parameter values of PYTHIA, respectively, while it is closer to the prediction of 5.48 from the PYTHIA parameter set from Ref. [8] and to the recent model predictions of 5.58 and 5.78 from Refs. [25] and [26]. The measured excess of the number of charged hadrons with respect to the event generators is independent of η and concentrated in the $p_T < 1$ GeV/c range. These differences indicate the need for a continued model development and simulation tuning. Work on updated event generators based on LHC data is currently underway.

Summary. Charged-hadron transverse-momentum and pseudorapidity distributions have been measured in proton-proton collisions at $\sqrt{s} = 7$ TeV. The numerical values of the data presented in this paper can be found in the HEPDATA database [32]. The combined result for the central pseudorapidity density, from three mutually consistent methods of measurement, is $dN_{ch}/d\eta|_{|\eta| < 0.5} = 5.78 \pm 0.01$ (stat.) ± 0.23 (syst.) for non-single-diffractive events. This value

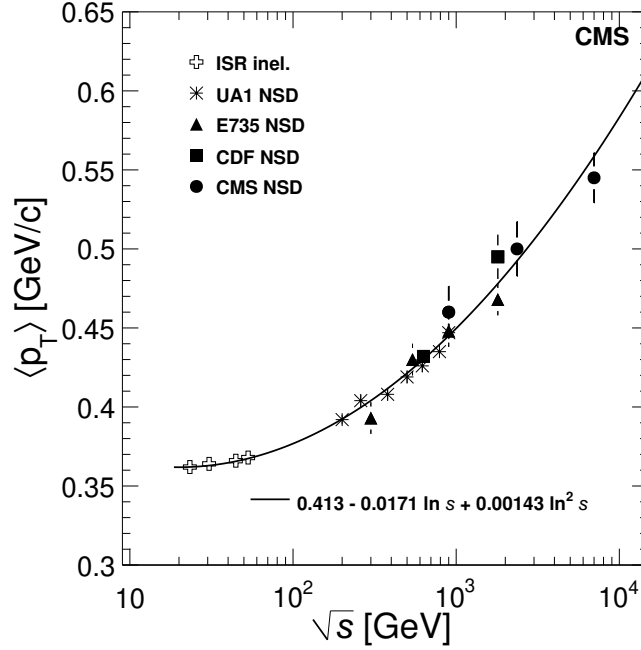


Figure 4: Average p_T of charged hadrons as a function of the centre-of-mass energy. The CMS measurements are for $|\eta| < 2.4$. Also shown are measurements from the ISR [19] (pp), E735 [20] ($p\bar{p}$), and CDF [21] ($p\bar{p}$) for $|\eta| < 0.5$, and from UA1 [18] ($p\bar{p}$) for $|\eta| < 2.5$. The solid line is a fit of the functional form $\langle p_T \rangle = 0.413 - 0.0171 \ln s + 0.00143 \ln^2 s$ to the data. The error bars on the CMS data include the systematic uncertainties.

is higher than most predictions and provides new information to constrain ongoing improvements of soft particle production models and event generators. The mean transverse momentum has been measured to be 0.545 ± 0.005 (stat.) ± 0.015 (syst.) GeV/c. These studies are the first steps in the exploration of particle production at the new centre-of-mass energy frontier, and contribute to the understanding of the dynamics in soft hadronic interactions.

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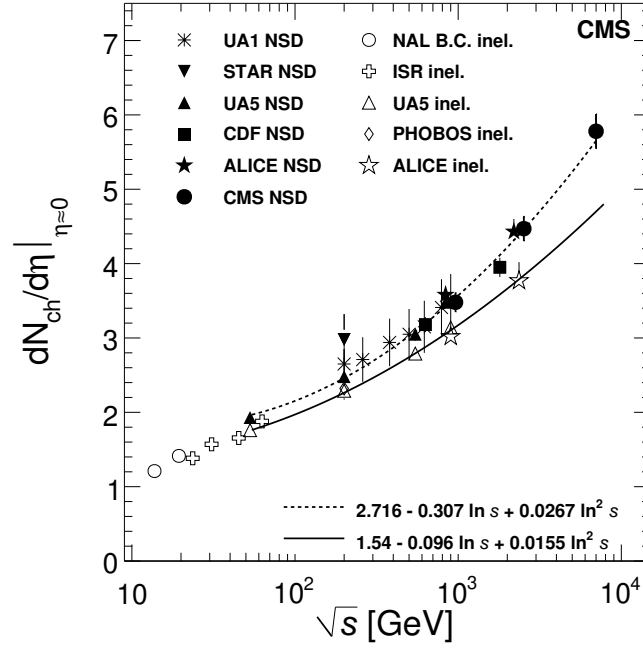


Figure 5: Average value of $dN_{\text{ch}}/d\eta$ in the central η region as a function of centre-of-mass energy in pp and $p\bar{p}$ collisions. Also shown are NSD and inelastic measurements from the NAL Bubble Chamber [27] ($p\bar{p}$), ISR [28] (pp), UA1 [18] ($p\bar{p}$), UA5 [15] ($p\bar{p}$), CDF [29] ($p\bar{p}$), STAR [30] (pp), PHOBOS [31] (pp), and ALICE [16] (pp). The curves are second-order polynomial fits for the inelastic (solid) and NSD event selections (dashed). The error bars include systematic uncertainties, when available. Data points at 0.9 and 2.36 TeV are slightly displaced horizontally for visibility.

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A The CMS Collaboration

Yerevan Physics Institute, Yerevan, Armenia

V. Khachatryan, A.M. Sirunyan, A. Tumasyan

Institut für Hochenergiephysik der OeAW, Wien, Austria

W. Adam, T. Bergauer, M. Dragicevic, J. Erö, C. Fabjan, M. Friedl, R. Frühwirth, V.M. Ghete, J. Hammer¹, S. Häseler, M. Hoch, N. Hörmann, J. Hrubec, M. Jeitler, G. Kasieczka, W. Kiesenhofer, M. Krammer, D. Liko, I. Mikulec, M. Pernicka, H. Rohringer, R. Schöfbeck, J. Strauss, A. Taurok, F. Teischinger, W. Waltenberger, G. Walzel, E. Widl, C.-E. Wulz

National Centre for Particle and High Energy Physics, Minsk, Belarus

V. Mossolov, N. Shumeiko, J. Suarez Gonzalez

Universiteit Antwerpen, Antwerpen, Belgium

L. Benucci, L. Ceard, E.A. De Wolf, M. Hashemi, X. Janssen, T. Maes, L. Mucibello, S. Ochesanu, B. Roland, R. Rougny, M. Selvaggi, H. Van Haevermaet, P. Van Mechelen, N. Van Remortel

Vrije Universiteit Brussel, Brussel, Belgium

V. Adler, S. Beauceron, S. Blyweert, J. D'Hondt, O. Devroede, A. Kalogeropoulos, J. Maes, M. Maes, S. Tavernier, W. Van Doninck, P. Van Mulders, I. Villella

Université Libre de Bruxelles, Bruxelles, Belgium

E.C. Chabert, O. Charaf, B. Clerbaux, G. De Lentdecker, V. Dero, A.P.R. Gay, G.H. Hammad, P.E. Marage, C. Vander Velde, P. Vanlaer, J. Wickens

Ghent University, Ghent, Belgium

S. Costantini, M. Grunewald, B. Klein, A. Marinov, D. Ryckbosch, F. Thyssen, M. Tytgat, L. Vanelderen, P. Verwilligen, S. Walsh, N. Zaganidis

Université Catholique de Louvain, Louvain-la-Neuve, Belgium

S. Basegmez, G. Bruno, J. Caudron, J. De Favereau De Jeneret, C. Delaere, P. Demin, D. Favart, A. Giammanco, G. Grégoire, J. Hollar, V. Lemaitre, O. Militaru, S. Oryn, D. Pagano, A. Pin, K. Piotrkowski¹, L. Quertenmont, N. Schul

Université de Mons, Mons, Belgium

N. Bely, T. Caebergs, E. Daubie

Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil

G.A. Alves, M.E. Pol, M.H.G. Souza

Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

W. Carvalho, E.M. Da Costa, D. De Jesus Damiao, C. De Oliveira Martins, S. Fonseca De Souza, L. Mundim, V. Oguri, A. Santoro, S.M. Silva Do Amaral, A. Sznajder, F. Torres Da Silva De Araujo

Instituto de Fisica Teorica, Universidade Estadual Paulista, Sao Paulo, Brazil

F.A. Dias, M.A.F. Dias, T.R. Fernandez Perez Tomei, E. M. Gregores², F. Marinho, S.F. Novaes, Sandra S. Padula

Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria

N. Darmanov¹, L. Dimitrov, V. Genchev¹, P. Iaydjiev, S. Piperov, S. Stoykova, G. Sultanov, R. Trayanov, I. Vankov

University of Sofia, Sofia, Bulgaria

M. Dyulendarova, R. Hadjiiska, V. Kozhuharov, L. Litov, E. Marinova, M. Mateev, B. Pavlov, P. Petkov

Institute of High Energy Physics, Beijing, China

J.G. Bian, G.M. Chen, H.S. Chen, C.H. Jiang, D. Liang, S. Liang, J. Wang, J. Wang, X. Wang, Z. Wang, M. Yang, J. Zang, Z. Zhang

State Key Lab. of Nucl. Phys. and Tech., Peking University, Beijing, China

Y. Ban, S. Guo, Z. Hu, Y. Mao, S.J. Qian, H. Teng, B. Zhu

Universidad de Los Andes, Bogota, Colombia

A. Cabrera, C.A. Carrillo Montoya, B. Gomez Moreno, A.A. Ocampo Rios, A.F. Osorio Oliveros, J.C. Sanabria

Technical University of Split, Split, Croatia

N. Godinovic, D. Lelas, K. Lelas, R. Plestina³, D. Polic, I. Puljak

University of Split, Split, Croatia

Z. Antunovic, M. Dzelalija

Institute Rudjer Boskovic, Zagreb, Croatia

V. Brigljevic, S. Duric, K. Kadija, S. Morovic

University of Cyprus, Nicosia, Cyprus

A. Attikis, R. Fereos, M. Galanti, J. Mousa, C. Nicolaou, A. Papadakis, F. Ptochos, P.A. Razis, H. Rykaczewski, D. Tsiakkouri, Z. Zinonos

Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt

M. Mahmoud⁴

National Institute of Chemical Physics and Biophysics, Tallinn, Estonia

A. Hektor, M. Kadastik, K. Kannike, M. Müntel, M. Raidal, L. Rebane

Department of Physics, University of Helsinki, Helsinki, Finland

V. Azzolini, P. Eerola

Helsinki Institute of Physics, Helsinki, Finland

S. Czellar, J. Härkönen, A. Heikkinen, V. Karimäki, R. Kinnunen, J. Klem, M.J. Kortelainen, T. Lampén, K. Lassila-Perini, S. Lehti, T. Lindén, P. Luukka, T. Mäenpää, E. Tuominen, J. Tuominiemi, E. Tuovinen, D. Ungaro, L. Wendland

Lappeenranta University of Technology, Lappeenranta, Finland

K. Banzuzi, A. Korpela, T. Tuuva

Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France

D. Sillou

DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France

M. Besancon, M. Dejardin, D. Denegri, J. Descamps, B. Fabbro, J.L. Faure, F. Ferri, S. Ganjour, F.X. Gentit, A. Givernaud, P. Gras, G. Hamel de Monchenault, P. Jarry, E. Locci, J. Malcles, M. Marionneau, L. Millischer, J. Rander, A. Rosowsky, D. Rousseau, M. Titov, P. Verrecchia

Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France

S. Baffioni, L. Bianchini, M. Bluj⁵, C. Broutin, P. Busson, C. Charlot, L. Dobrzynski, S. Elgammal,

R. Granier de Cassagnac, M. Haguenauer, A. Kalinowski, P. Miné, P. Paganini, D. Sabes, Y. Sirois, C. Thiebaux, A. Zabi

Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France

J.-L. Agram⁶, A. Besson, D. Bloch, D. Bodin, J.-M. Brom, M. Cardaci, E. Conte⁶, F. Drouhin⁶, C. Ferro, J.-C. Fontaine⁶, D. Gelé, U. Goerlach, S. Greder, P. Juillot, M. Karim⁶, A.-C. Le Bihan, Y. Mikami, J. Speck, P. Van Hove

Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules (IN2P3), Villeurbanne, France

F. Fassi, D. Mercier

Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France

C. Baty, N. Beaupere, M. Bedjidian, O. Bondu, G. Boudoul, D. Boumediene, H. Brun, N. Chanon, R. Chierici, D. Contardo, P. Depasse, H. El Mamouni, J. Fay, S. Gascon, B. Ille, T. Kurca, T. Le Grand, M. Lethuillier, L. Mirabito, S. Perries, S. Tosi, Y. Tschudi, P. Verdier, H. Xiao

E. Andronikashvili Institute of Physics, Academy of Science, Tbilisi, Georgia

V. Roinishvili

RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany

G. Anagnostou, M. Edelhoff, L. Feld, N. Heracleous, O. Hindrichs, R. Jussen, K. Klein, J. Merz, N. Mohr, A. Ostapchuk, A. Perieanu, F. Raupach, J. Sammet, S. Schael, D. Sprenger, H. Weber, M. Weber, B. Wittmer

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

O. Actis, M. Ata, W. Bender, P. Biallass, M. Erdmann, J. Frangenheim, T. Hebbeker, A. Hinzmann, K. Hoepfner, C. Hof, M. Kirsch, T. Klimkovich, P. Kreuzer¹, D. Lanske[†], C. Magass, M. Merschmeyer, A. Meyer, P. Papacz, H. Pieta, H. Reithler, S.A. Schmitz, L. Sonnenschein, M. Sowa, J. Steggemann, D. Teyssier, C. Zeidler

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

M. Bontenackels, M. Davids, M. Duda, G. Flügge, H. Geenen, M. Giffels, W. Haj Ahmad, D. Heydhausen, T. Kress, Y. Kuessel, A. Linn, A. Nowack, L. Perchalla, O. Pooth, P. Sauerland, A. Stahl, M. Thomas, D. Tornier, M.H. Zoeller

Deutsches Elektronen-Synchrotron, Hamburg, Germany

M. Aldaya Martin, W. Behrenhoff, U. Behrens, M. Bergholz, K. Borras, A. Campbell, E. Castro, D. Dammann, G. Eckerlin, A. Flossdorf, G. Flucke, A. Geiser, J. Hauk, H. Jung, M. Kasemann, I. Katkov, C. Kleinwort, H. Kluge, A. Knutsson, E. Kuznetsova, W. Lange, W. Lohmann, R. Mankel, M. Marienfeld, I.-A. Melzer-Pellmann, A.B. Meyer, J. Mnich, A. Mussgiller, J. Olzem, A. Parenti, A. Raspereza, R. Schmidt, T. Schoerner-Sadenius, N. Sen, M. Stein, J. Tomaszewska, D. Volyansky, C. Wissing

University of Hamburg, Hamburg, Germany

C. Autermann, J. Draeger, D. Eckstein, H. Enderle, U. Gebbert, K. Kaschube, G. Kaussen, R. Klanner, B. Mura, S. Naumann-Emme, F. Nowak, C. Sander, H. Schettler, P. Schleper, M. Schröder, T. Schum, J. Schwandt, H. Stadie, G. Steinbrück, J. Thomsen, R. Wolf

Institut für Experimentelle Kernphysik, Karlsruhe, Germany

J. Bauer, V. Buege, A. Cakir, T. Chwalek, D. Daeuwel, W. De Boer, A. Dierlamm, G. Dirkes,

M. Feindt, J. Gruschke, C. Hackstein, F. Hartmann, M. Heinrich, H. Held, K.H. Hoffmann, S. Honc, T. Kuhr, D. Martschei, S. Mueller, Th. Müller, M. Niegel, O. Oberst, A. Oehler, J. Ott, T. Peiffer, D. Piparo, G. Quast, K. Rabbertz, F. Ratnikov, M. Renz, A. Sabellek, C. Saout¹, A. Scheurer, P. Schieferdecker, F.-P. Schilling, G. Schott, H.J. Simonis, F.M. Stober, D. Troendle, J. Wagner-Kuhr, M. Zeise, V. Zhukov⁷, E.B. Ziebarth

Institute of Nuclear Physics "Demokritos", Aghia Paraskevi, Greece

G. Daskalakis, T. Gerasis, A. Kyriakis, D. Loukas, I. Manolakos, A. Markou, C. Markou, C. Mavrommatis, E. Petrakou

University of Athens, Athens, Greece

L. Gouskos, P. Katsas, A. Panagiotou¹

University of Ioánnina, Ioánnina, Greece

I. Evangelou, P. Kokkas, N. Manthos, I. Papadopoulos, V. Patras, F.A. Triantis

KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary

A. Aranyi, G. Bencze, L. Boldizsar, G. Debreczeni, C. Hajdu¹, D. Horvath⁸, A. Kapusi, K. Krajczar⁹, A. Laszlo, F. Sikler, G. Vesztergombi⁹

Institute of Nuclear Research ATOMKI, Debrecen, Hungary

N. Beni, J. Molnar, J. Palinkas, Z. Szillasi¹, V. Veszpremi

University of Debrecen, Debrecen, Hungary

P. Raics, Z.L. Trocsanyi, B. Ujvari

Panjab University, Chandigarh, India

S. Bansal, S.B. Beri, V. Bhatnagar, M. Jindal, M. Kaur, J.M. Kohli, M.Z. Mehta, N. Nishu, L.K. Saini, A. Sharma, R. Sharma, A.P. Singh, J.B. Singh, S.P. Singh

University of Delhi, Delhi, India

S. Ahuja, S. Bhattacharya¹⁰, S. Chauhan, B.C. Choudhary, P. Gupta, S. Jain, S. Jain, A. Kumar, K. Ranjan, R.K. Shivpuri

Bhabha Atomic Research Centre, Mumbai, India

R.K. Choudhury, D. Dutta, S. Kailas, S.K. Kataria, A.K. Mohanty, L.M. Pant, P. Shukla, P. Suggiseti

Tata Institute of Fundamental Research - EHEP, Mumbai, India

T. Aziz, M. Guchait¹¹, A. Gurtu, M. Maity¹², D. Majumder, G. Majumder, K. Mazumdar, G.B. Mohanty, A. Saha, K. Sudhakar, N. Wickramage

Tata Institute of Fundamental Research - HECR, Mumbai, India

S. Banerjee, S. Dugad, N.K. Mondal

Institute for Studies in Theoretical Physics & Mathematics (IPM), Tehran, Iran

H. Arfaei, H. Bakhshiansohi, A. Fahim, A. Jafari, M. Mohammadi Najafabadi, S. Paktinat Mehdiabadi, B. Safarzadeh, M. Zeinali

INFN Sezione di Bari ^a, Università di Bari ^b, Politecnico di Bari ^c, Bari, Italy

M. Abbrescia^{a,b}, L. Barbone^a, A. Colaleo^a, D. Creanza^{a,c}, N. De Filippis^a, M. De Palma^{a,b}, A. Dimitrov^a, F. Fedele^a, L. Fiore^a, G. Iaselli^{a,c}, L. Lusito^{a,b,1}, G. Maggi^{a,c}, M. Maggi^a, N. Manna^{a,b}, B. Marangelli^{a,b}, S. My^{a,c}, S. Nuzzo^{a,b}, G.A. Pierro^a, A. Pompili^{a,b}, G. Pugliese^{a,c}, F. Romano^{a,c}, G. Roselli^{a,b}, G. Selvaggi^{a,b}, L. Silvestris^a, R. Trentadue^a, S. Tupputi^{a,b}, G. Zito^a

INFN Sezione di Bologna ^a, Università di Bologna ^b, Bologna, Italy

G. Abbiendi^a, A.C. Benvenuti^a, D. Bonacorsi^a, S. Braibant-Giacomelli^{a,b}, A. Castro^{a,b}, F.R. Cavallo^a, G. Codispoti^{a,b}, G.M. Dallavalle^{a,1}, F. Fabbri^a, A. Fanfani^{a,b}, D. Fasanella^a, P. Giacomelli^a, M. Giunta^{a,1}, C. Grandi^a, S. Marcellini^a, G. Masetti^{a,b}, A. Montanari^a, F.L. Navarria^{a,b}, F. Odorici^a, A. Perrotta^a, A.M. Rossi^{a,b}, T. Rovelli^{a,b}, G. Siroli^{a,b}, R. Travaglini^{a,b}

INFN Sezione di Catania ^a, Università di Catania ^b, Catania, Italy

S. Albergo^{a,b}, G. Cappello^{a,b}, M. Chiorboli^{a,b}, S. Costa^{a,b}, A. Tricomi^{a,b}, C. Tuve^a

INFN Sezione di Firenze ^a, Università di Firenze ^b, Firenze, Italy

G. Barbagli^a, G. Broccolo^{a,b}, V. Ciulli^{a,b}, C. Civinini^a, R. D'Alessandro^{a,b}, E. Focardi^{a,b}, S. Frosali^{a,b}, E. Gallo^a, C. Genta^{a,b}, P. Lenzi^{a,b,1}, M. Meschini^a, S. Paoletti^a, G. Sguazzoni^a, A. Tropiano^a

INFN Laboratori Nazionali di Frascati, Frascati, Italy

L. Benussi, S. Bianco, S. Colafranceschi¹³, F. Fabbri, D. Piccolo

INFN Sezione di Genova, Genova, Italy

P. Fabbriatore, R. Musenich

INFN Sezione di Milano-Bicocca ^a, Università di Milano-Bicocca ^b, Milano, Italy

A. Benaglia^{a,b}, G.B. Cerati^{a,b,1}, F. De Guio^{a,b}, L. Di Matteo^{a,b}, A. Ghezzi^{a,b,1}, P. Govoni^{a,b}, M. Malberti^{a,b,1}, S. Malvezzi^a, A. Martelli^{a,b,3}, A. Massironi^{a,b}, D. Menasce^a, V. Miccio^{a,b}, L. Moroni^a, P. Negri^{a,b}, M. Paganoni^{a,b}, D. Pedrini^a, S. Ragazzi^{a,b}, N. Redaelli^a, S. Sala^a, R. Salerno^{a,b}, T. Tabarelli de Fatis^{a,b}, V. Tancini^{a,b}, S. Taroni^{a,b}

INFN Sezione di Napoli ^a, Università di Napoli "Federico II" ^b, Napoli, Italy

S. Buontempo^a, A. Cimmino^{a,b}, A. De Cosa^{a,b,1}, M. De Gruttola^{a,b,1}, F. Fabozzi^{a,14}, A.O.M. Iorio^a, L. Lista^a, P. Noli^{a,b}, P. Paolucci^a

INFN Sezione di Padova ^a, Università di Padova ^b, Università di Trento (Trento) ^c, Padova, Italy

P. Azzi^a, N. Bacchetta^a, P. Bellan^{a,b,1}, M. Bellato^a, M. Biasotto^{a,15}, D. Bisello^{a,b}, R. Carlin^{a,b}, P. Checchia^a, M. De Mattia^{a,b}, T. Dorigo^a, F. Fanzago^a, F. Gasparini^{a,b}, P. Giubilato^{a,b}, A. Gresele^{a,c}, S. Lacaprara^a, I. Lazzizzera^{a,c}, M. Margoni^{a,b}, G. Maron^{a,15}, A.T. Meneguzzo^{a,b}, M. Nespolo^a, L. Perrozzi^a, N. Pozzobon^{a,b}, P. Ronchese^{a,b}, F. Simonetto^{a,b}, E. Torassa^a, M. Tosi^{a,b}, A. Triossi^a, S. Vanini^{a,b}, G. Zumerle^{a,b}

INFN Sezione di Pavia ^a, Università di Pavia ^b, Pavia, Italy

P. Baesso^{a,b}, U. Berzano^a, C. Riccardi^{a,b}, P. Torre^{a,b}, P. Vitulo^{a,b}, C. Viviani^{a,b}

INFN Sezione di Perugia ^a, Università di Perugia ^b, Perugia, Italy

M. Biasini^{a,b}, G.M. Bilei^a, B. Caponeri^{a,b}, L. Fanò^a, P. Lariccia^{a,b}, A. Lucaroni^{a,b}, G. Mantovani^{a,b}, M. Menichelli^a, A. Nappi^{a,b}, A. Santocchia^{a,b}, L. Servoli^a, M. Valdata^a, R. Volpe^{a,b,1}

INFN Sezione di Pisa ^a, Università di Pisa ^b, Scuola Normale Superiore di Pisa ^c, Pisa, Italy

P. Azzurri^{a,c}, G. Bagliesi^a, J. Bernardini^{a,b,1}, T. Boccali^a, R. Castaldi^a, R.T. Dagnolo^{a,c}, R. Dell'Orso^a, F. Fiori^{a,b}, L. Foà^{a,c}, A. Giassi^a, A. Kraan^a, F. Ligabue^{a,c}, T. Lomtadze^a, L. Martini^a, A. Messineo^{a,b}, F. Palla^a, F. Palmonari^a, G. Segneri^a, A.T. Serban^a, P. Spagnolo^{a,1}, R. Tenchini^{a,1}, G. Tonelli^{a,b,1}, A. Venturi^a, P.G. Verdini^a

INFN Sezione di Roma ^a, Università di Roma "La Sapienza" ^b, Roma, Italy

L. Barone^{a,b}, F. Cavallari^{a,1}, D. Del Re^{a,b}, E. Di Marco^{a,b}, M. Diemoz^a, D. Franci^{a,b}, M. Grassi^a, E. Longo^{a,b}, G. Organtini^{a,b}, A. Palma^{a,b}, F. Pandolfi^{a,b}, R. Paramatti^{a,1}, S. Rahatlou^{a,b,1}

INFN Sezione di Torino ^a, Università di Torino ^b, Università del Piemonte Orientale (Novara) ^c, Torino, Italy

N. Amapane^{a,b}, R. Arcidiacono^{a,b}, S. Argiro^{a,b}, M. Arneodo^{a,c}, C. Biino^a, C. Botta^{a,b}, N. Cartiglia^a, R. Castello^{a,b}, M. Costa^{a,b}, N. Demaria^a, A. Graziano^{a,b}, C. Mariotti^a, M. Marone^{a,b}, S. Maselli^a, E. Migliore^{a,b}, G. Mila^{a,b}, V. Monaco^{a,b}, M. Musich^{a,b}, M.M. Obertino^{a,c}, N. Pastrone^a, M. Pelliccioni^{a,b,1}, A. Romero^{a,b}, M. Ruspa^{a,c}, R. Sacchi^{a,b}, A. Solano^{a,b}, A. Staiano^a, D. Trocino^{a,b}, A. Vilela Pereira^{a,b,1}

INFN Sezione di Trieste ^a, Università di Trieste ^b, Trieste, Italy

F. Ambroglini^{a,b}, S. Belforte^a, F. Cossutti^a, G. Della Ricca^{a,b}, B. Gobbo^a, D. Montanino^a, A. Penzo^a

Kyungpook National University, Daegu, Korea

S. Chang, J. Chung, D.H. Kim, G.N. Kim, J.E. Kim, D.J. Kong, H. Park, D.C. Son

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea

Zero Kim, J.Y. Kim, S. Song

Korea University, Seoul, Korea

B. Hong, H. Kim, J.H. Kim, T.J. Kim, K.S. Lee, D.H. Moon, S.K. Park, H.B. Rhee, K.S. Sim

University of Seoul, Seoul, Korea

M. Choi, S. Kang, H. Kim, C. Park, I.C. Park, S. Park

Sungkyunkwan University, Suwon, Korea

S. Choi, Y. Choi, Y.K. Choi, J. Goh, J. Lee, S. Lee, H. Seo, I. Yu

Vilnius University, Vilnius, Lithuania

M. Janulis, D. Martisiute, P. Petrov, T. Sabonis

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

H. Castilla Valdez¹, E. De La Cruz Burelo, R. Lopez-Fernandez, A. Sánchez Hernández, L.M. Villaseñor-Cendejas

Universidad Iberoamericana, Mexico City, Mexico

S. Carrillo Moreno

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

H.A. Salazar Ibarguen

Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico

E. Casimiro Linares, A. Morelos Pineda, M.A. Reyes-Santos

University of Auckland, Auckland, New Zealand

P. Allfrey, D. Krofcheck, J. Tam

University of Canterbury, Christchurch, New Zealand

T. Aumeyr, P.H. Butler, T. Signal, J.C. Williams

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

M. Ahmad, I. Ahmed, M.I. Asghar, H.R. Hoorani, W.A. Khan, T. Khurshid, S. Qazi

Institute of Experimental Physics, Warsaw, Poland

M. Cwiok, W. Dominik, K. Doroba, M. Konecki, J. Krolikowski

Soltan Institute for Nuclear Studies, Warsaw, Poland

T. Frueboes, R. Gokieli, M. Górski, M. Kazana, K. Nawrocki, M. Szleper, G. Wrochna, P. Zalewski

Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

N. Almeida, A. David, P. Faccioli, P.G. Ferreira Parracho, M. Gallinaro, G. Mini, P. Musella, A. Nayak, L. Raposo, P.Q. Ribeiro, J. Seixas, P. Silva, D. Soares, J. Varela¹, H.K. Wöhri

Joint Institute for Nuclear Research, Dubna, Russia

I. Altsybeev, I. Belotelov, P. Bunin, M. Finger, M. Finger Jr., I. Golutvin, A. Kamenev, V. Karjavin, G. Kozlov, A. Lanev, P. Moisenz, V. Palichik, V. Perelygin, S. Shmatov, V. Smirnov, A. Volodko, A. Zarubin

Petersburg Nuclear Physics Institute, Gatchina (St Petersburg), Russia

N. Bondar, V. Golovtsov, Y. Ivanov, V. Kim, P. Levchenko, I. Smirnov, V. Sulimov, L. Uvarov, S. Vavilov, A. Vorobyev

Institute for Nuclear Research, Moscow, Russia

Yu. Andreev, S. Gninenko, N. Golubev, M. Kirsanov, N. Krasnikov, V. Matveev, A. Pashenkov, A. Toropin, S. Troitsky

Institute for Theoretical and Experimental Physics, Moscow, Russia

V. Epshteyn, V. Gavrilov, N. Ilina, V. Kaftanov[†], M. Kossov¹, A. Krokhotin, S. Kuleshov, A. Oulianov, G. Safronov, S. Semenov, I. Shreyber, V. Stolin, E. Vlasov, A. Zhokin

Moscow State University, Moscow, Russia

E. Boos, M. Dubinin¹⁶, L. Dudko, A. Ershov, A. Gribushin, O. Kodolova, I. Lokhtin, S. Obraztsov, S. Petrushanko, L. Sarycheva, V. Savrin, A. Snigirev

P.N. Lebedev Physical Institute, Moscow, Russia

V. Andreev, I. Dremin, M. Kirakosyan, S.V. Rusakov, A. Vinogradov

State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia

I. Azhgirey, S. Bitioukov, K. Datsko, V. Grishin¹, V. Kachanov, D. Konstantinov, V. Krychkine, V. Petrov, R. Ryutin, S. Slabospitsky, A. Sobol, A. Sytine, L. Tourtchanovitch, S. Troshin, N. Tyurin, A. Uzunian, A. Volkov

Vinca Institute of Nuclear Sciences, Belgrade, Serbia

P. Adzic¹⁷, M. Djordjevic, D. Krpic¹⁷, D. Maletic, J. Milosevic, J. Puzovic

Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain

M. Aguilar-Benitez, J. Alcaraz Maestre, P. Arce, C. Battilana, E. Calvo, M. Cepeda, M. Cerrada, M. Chamizo Llatas, N. Colino, B. De La Cruz, C. Diez Pardos, C. Fernandez Bedoya, J.P. Fernández Ramos, A. Ferrando, J. Flix, M.C. Fouz, P. Garcia-Abia, O. Gonzalez Lopez, S. Goy Lopez, J.M. Hernandez, M.I. Josa, G. Merino, J. Puerta Pelayo, I. Redondo, L. Romero, J. Santaolalla, C. Willmott

Universidad Autónoma de Madrid, Madrid, Spain

C. Albajar, J.F. de Trocóniz

Universidad de Oviedo, Oviedo, Spain

J. Cuevas, J. Fernandez Menendez, I. Gonzalez Caballero, L. Lloret Iglesias, J.M. Vizan Garcia

Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain

I.J. Cabrillo, A. Calderon, S.H. Chuang, I. Diaz Merino, C. Diez Gonzalez, J. Duarte Campderros, M. Fernandez, G. Gomez, J. Gonzalez Sanchez, R. Gonzalez Suarez, C. Jorda, P. Lobelle Pardo, A. Lopez Virto, J. Marco, R. Marco, C. Martinez Rivero, P. Martinez Ruiz del Arbol, F. Matorras, T. Rodrigo, A. Ruiz Jimeno, L. Scodellaro, M. Sobron Sanudo, I. Vila, R. Vilar Cortabitarte

CERN, European Organization for Nuclear Research, Geneva, Switzerland

D. Abbaneo, E. Auffray, P. Baillon, A.H. Ball, D. Barney, F. Beaudette³, A.J. Bell¹⁸, R. Bellan, D. Benedetti, C. Bernet³, W. Bialas, P. Bloch, A. Bocci, S. Bolognesi, H. Breuker, G. Brona, K. Bunkowski, T. Camporesi, E. Cano, A. Cattai, G. Cerminara, T. Christiansen, J.A. Coarasa Perez, R. Covarelli, B. Curé, T. Dahms, A. De Roeck, A. Elliott-Peisert, W. Funk, A. Gaddi, S. Gennai, H. Gerwig, D. Gigi, K. Gill, D. Giordano, F. Glege, R. Gomez-Reino Garrido, S. Gowdy, L. Guiducci, M. Hansen, C. Hartl, J. Harvey, B. Hegner, C. Henderson, H.F. Hoffmann, A. Honma, V. Innocente, P. Janot, P. Lecoq, C. Leonidopoulos, C. Lourenço, A. Macpherson, T. Mäki, L. Malgeri, M. Mannelli, L. Masetti, G. Mavromanolakis, F. Meijers, S. Mersi, E. Meschi, R. Moser, M.U. Mozer, M. Mulders, E. Nesvold¹, L. Orsini, E. Perez, A. Petrilli, A. Pfeiffer, M. Pierini, M. Pimiä, A. Racz, G. Rolandi¹⁹, C. Rovelli²⁰, M. Rovere, V. Ryjov, H. Sakulin, C. Schäfer, C. Schwick, I. Segoni, A. Sharma, P. Siegrist, M. Simon, P. Sphicas²¹, D. Spiga, M. Spiropulu¹⁶, F. Stöckli, P. Traczyk, P. Tropea, A. Tsirou, G.I. Veres⁹, P. Vichoudis, M. Voutilainen, W.D. Zeuner

Paul Scherrer Institut, Villigen, Switzerland

W. Bertl, K. Deiters, W. Erdmann, K. Gabathuler, R. Horisberger, Q. Ingram, H.C. Kaestli, S. König, D. Kotlinski, U. Langenegger, F. Meier, D. Renker, T. Rohe, J. Sibille²², A. Starodumov²³

Institute for Particle Physics, ETH Zurich, Zurich, Switzerland

L. Caminada²⁴, Z. Chen, S. Cittolin, G. Dissertori, M. Dittmar, J. Eugster, K. Freudenreich, C. Grab, A. Hervé, W. Hintz, P. Lecomte, W. Lustermann, C. Marchica²⁴, P. Meridiani, P. Milenovic²⁵, F. Moortgat, A. Nardulli, F. Nessi-Tedaldi, L. Pape, F. Pauss, T. Punz, A. Rizzi, F.J. Ronga, L. Sala, A.K. Sanchez, M.-C. Sawley, D. Schinzel, V. Sordini, B. Stieger, L. Tauscher[†], A. Thea, K. Theofilatos, D. Treille, M. Weber, L. Wehrli, J. Weng

Universität Zürich, Zurich, Switzerland

C. Amsler, V. Chiochia, S. De Visscher, M. Ivova Rikova, B. Millan Mejias, C. Regenfus, P. Robmann, T. Rommerskirchen, A. Schmidt, D. Tsirigkas, L. Wilke

National Central University, Chung-Li, Taiwan

Y.H. Chang, K.H. Chen, W.T. Chen, A. Go, C.M. Kuo, S.W. Li, W. Lin, M.H. Liu, Y.J. Lu, J.H. Wu, S.S. Yu

National Taiwan University (NTU), Taipei, Taiwan

P. Bartalini, P. Chang, Y.H. Chang, Y.W. Chang, Y. Chao, K.F. Chen, W.-S. Hou, Y. Hsiung, K.Y. Kao, Y.J. Lei, S.W. Lin, R.-S. Lu, J.G. Shiu, Y.M. Tzeng, K. Ueno, C.C. Wang, M. Wang, J.T. Wei

Cukurova University, Adana, Turkey

A. Adiguzel, A. Ayhan, M.N. Bakirci, S. Cerci²⁶, Z. Demir, C. Dozen, I. Dumanoglu, E. Eskut, S. Girgis, G. Gökbulut, Y. Güler, E. Gurpinar, I. Hos, E.E. Kangal, T. Karaman, A. Kayis Topaksu, A. Nart, G. Önengüt, K. Ozdemir, S. Ozturk, A. Polatöz, O. Sahin, O. Sengul, K. Sogut²⁷, B. Tali, H. Topakli, D. Uzun, L.N. Vergili, M. Vergili, C. Zorbilmez

Middle East Technical University, Physics Department, Ankara, Turkey

I.V. Akin, T. Aliev, S. Bilmis, M. Deniz, H. Gamsizkan, A.M. Guler, K. Ocalan, A. Ozpineci, M. Serin, R. Sever, U.E. Surat, M. Zeyrek

Bogaziçi University, Department of Physics, Istanbul, Turkey

M. Deliomeroglu, D. Demir²⁸, E. Gülmez, A. Halu, B. Isildak, M. Kaya²⁹, O. Kaya²⁹, M. Özbek, S. Ozkorucuklu³⁰, N. Sonmez³¹

National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine

L. Levchuk

University of Bristol, Bristol, United Kingdom

P. Bell, F. Bostock, J.J. Brooke, T.L. Cheng, D. Cussans, R. Frazier, J. Goldstein, M. Hansen, G.P. Heath, H.F. Heath, C. Hill, B. Huckvale, J. Jackson, L. Kreczko, C.K. Mackay, S. Metson, D.M. Newbold³², K. Nirunpong, V.J. Smith, S. Ward

Rutherford Appleton Laboratory, Didcot, United Kingdom

L. Basso, K.W. Bell, A. Belyaev, C. Brew, R.M. Brown, B. Camanzi, D.J.A. Cockerill, J.A. Coughlan, K. Harder, S. Harper, B.W. Kennedy, E. Olaiya, B.C. Radburn-Smith, C.H. Shepherd-Themistocleous, I.R. Tomalin, W.J. Womersley, S.D. Worm

Imperial College, University of London, London, United Kingdom

R. Bainbridge, G. Ball, J. Ballin, R. Beuselinck, O. Buchmuller, D. Colling, N. Cripps, M. Cutajar, G. Davies, M. Della Negra, C. Foudas, J. Fulcher, D. Futyan, A. Guneratne Bryer, G. Hall, Z. Hatherell, J. Hays, G. Iles, G. Karapostoli, L. Lyons, A.-M. Magnan, J. Marrouche, R. Nandi, J. Nash, A. Nikitenko²³, A. Papageorgiou, M. Pesaresi, K. Petridis, M. Pioppi³³, D.M. Raymond, N. Rompotis, A. Rose, M.J. Ryan, C. Seez, P. Sharp, A. Sparrow, M. Stoye, A. Tapper, S. Tournier, M. Vazquez Acosta, T. Virdee¹, S. Wakefield, D. Wardrope, T. Whyntie

Brunel University, Uxbridge, United Kingdom

M. Barrett, M. Chadwick, J.E. Cole, P.R. Hobson, A. Khan, P. Kyberd, D. Leslie, I.D. Reid, L. Teodorescu

Boston University, Boston, USA

T. Bose, A. Clough, A. Heister, J. St. John, P. Lawson, D. Lazic, J. Rohlf, L. Sulak

Brown University, Providence, USA

J. Andrea, A. Avetisyan, S. Bhattacharya, J.P. Chou, D. Cutts, S. Esen, U. Heintz, S. Jabeen, G. Kukartsev, G. Landsberg, M. Narain, D. Nguyen, T. Speer, K.V. Tsang

University of California, Davis, Davis, USA

M.A. Borgia, R. Breedon, M. Calderon De La Barca Sanchez, D. Cebra, M. Chertok, J. Conway, P.T. Cox, J. Dolen, R. Erbacher, E. Friis, W. Ko, A. Kopecky, R. Lander, H. Liu, S. Maruyama, T. Miceli, M. Nikolic, D. Pellett, J. Robles, T. Schwarz, M. Searle, J. Smith, M. Squires, M. Tripathi, R. Vasquez Sierra, C. Veelken

University of California, Los Angeles, Los Angeles, USA

V. Andreev, K. Arisaka, D. Cline, R. Cousins, A. Deisher, S. Erhan¹, C. Farrell, M. Felcini, J. Hauser, M. Ignatenko, C. Jarvis, C. Plager, G. Rakness, P. Schlein[†], J. Tucker, V. Valuev, R. Wallny

University of California, Riverside, Riverside, USA

J. Babb, R. Clare, J. Ellison, J.W. Gary, G. Hanson, G.Y. Jeng, S.C. Kao, F. Liu, H. Liu, A. Luthra, H. Nguyen, G. Pasztor³⁴, A. Satpathy, B.C. Shen[†], R. Stringer, J. Sturdy, S. Sumowidagdo, R. Wilken, S. Wimpenny

University of California, San Diego, La Jolla, USA

W. Andrews, J.G. Branson, E. Dusinger, D. Evans, F. Golf, A. Holzner, R. Kelley, M. Lebourgeois, J. Letts, B. Mangano, J. Muelmenstaedt, S. Padhi, C. Palmer, G. Petrucciani, H. Pi, M. Pieri, R. Ranieri, M. Sani, V. Sharma¹, S. Simon, Y. Tu, A. Vartak, F. Würthwein, A. Yagil

University of California, Santa Barbara, Santa Barbara, USA

D. Barge, M. Blume, C. Campagnari, M. D'Alfonso, T. Danielson, J. Garberson, J. Incandela, C. Justus, P. Kalavase, S.A. Koay, D. Kovalskyi, V. Krutelyov, J. Lamb, S. Lowette, V. Pavlunin, F. Rebassoo, J. Ribnik, J. Richman, R. Rossin, D. Stuart, W. To, J.R. Vlimant, M. Witherell

California Institute of Technology, Pasadena, USA

A. Bornheim, J. Bunn, M. Gataullin, D. Kcira, V. Litvine, Y. Ma, H.B. Newman, C. Rogan, K. Shin, V. Timciuc, J. Veverka, R. Wilkinson, Y. Yang, R.Y. Zhu

Carnegie Mellon University, Pittsburgh, USA

B. Akgun, R. Carroll, T. Ferguson, D.W. Jang, S.Y. Jun, M. Paulini, J. Russ, N. Terentyev, H. Vogel, I. Vorobiev

University of Colorado at Boulder, Boulder, USA

J.P. Cumalat, M.E. Dinardo, B.R. Drell, W.T. Ford, B. Heyburn, E. Luigi Lopez, U. Nauenberg, J.G. Smith, K. Stenson, K.A. Ulmer, S.R. Wagner, S.L. Zang

Cornell University, Ithaca, USA

L. Agostino, J. Alexander, F. Blekman, A. Chatterjee, S. Das, N. Eggert, L.J. Fields, L.K. Gibbons, B. Heltsley, W. Hopkins, A. Khukhunaishvili, B. Kreis, V. Kuznetsov, G. Nicolas Kaufman, J.R. Patterson, D. Puigh, D. Riley, A. Ryd, X. Shi, W. Sun, W.D. Teo, J. Thom, J. Thompson, J. Vaughan, Y. Weng, P. Wittich

Fairfield University, Fairfield, USA

A. Biselli, G. Cirino, D. Winn

Fermi National Accelerator Laboratory, Batavia, USA

S. Abdullin, M. Albrow, J. Anderson, G. Apollinari, M. Atac, J.A. Bakken, S. Banerjee, L.A.T. Bauerdick, A. Beretvas, J. Berryhill, P.C. Bhat, I. Bloch, F. Borchering, K. Burkett, J.N. Butler, V. Chetluru, H.W.K. Cheung, F. Chlebana, S. Cihangir, M. Demarteau, D.P. Eartly, V.D. Elvira, I. Fisk, J. Freeman, Y. Gao, E. Gottschalk, D. Green, O. Gutsche, A. Hahn, J. Hanlon, R.M. Harris, E. James, H. Jensen, M. Johnson, U. Joshi, R. Khatiwada, B. Kilminster, B. Klima, K. Kousouris, S. Kunori, S. Kwan, P. Limon, R. Lipton, J. Lykken, K. Maeshima, J.M. Marraffino, D. Mason, P. McBride, T. McCauley, T. Miao, K. Mishra, S. Mrenna, Y. Musienko³⁵, C. Newman-Holmes, V. O'Dell, S. Popescu, R. Pordes, O. Prokofyev, N. Saoulidou, E. Sexton-Kennedy, S. Sharma, R.P. Smith[†], A. Soha, W.J. Spalding, L. Spiegel, P. Tan, L. Taylor, S. Tkaczyk, L. Uplegger, E.W. Vaandering, R. Vidal, J. Whitmore, W. Wu, F. Yumiceva, J.C. Yun

University of Florida, Gainesville, USA

D. Acosta, P. Avery, D. Bourilkov, M. Chen, G.P. Di Giovanni, D. Dobur, A. Drozdetskiy, R.D. Field, Y. Fu, I.K. Furic, J. Gartner, B. Kim, S. Klimentenko, J. Konigsberg, A. Korytov, K. Kotov, A. Kropivnitskaya, T. Kypreos, K. Matchev, G. Mitselmakher, Y. Pakhotin, J. Piedra Gomez, C. Prescott, R. Remington, M. Schmitt, B. Scurlock, P. Sellers, D. Wang, J. Yelton, M. Zakaria

Florida International University, Miami, USA

C. Ceron, V. Gaultney, L. Kramer, L.M. Lebolo, S. Linn, P. Markowitz, G. Martinez, D. Mesa, J.L. Rodriguez

Florida State University, Tallahassee, USA

T. Adams, A. Askew, J. Chen, B. Diamond, S.V. Gleyzer, J. Haas, S. Hagopian, V. Hagopian, M. Jenkins, K.F. Johnson, H. Prosper, S. Sekmen, V. Veeraraghavan

Florida Institute of Technology, Melbourne, USA

M.M. Baarmand, S. Guragain, M. Hohlmann, H. Kalakhety, H. Mermerkaya, R. Ralich, I. Vodopiyanov

University of Illinois at Chicago (UIC), Chicago, USA

M.R. Adams, I.M. Anghel, L. Apanasevich, V.E. Bazterra, R.R. Betts, J. Callner, R. Cavanaugh, C. Dragoiu, E.J. Garcia-Solis, C.E. Gerber, D.J. Hofman, S. Khalatian, F. Lacroix, E. Shabalina, A. Smoron, D. Strom, N. Varelas

The University of Iowa, Iowa City, USA

U. Akgun, E.A. Albayrak, B. Bilki, K. Cankocak³⁶, W. Clarida, F. Duru, C.K. Lae, E. McCliment, J.-P. Merlo, A. Mestvirishvili, A. Moeller, J. Nachtman, C.R. Newsom, E. Norbeck, J. Olson, Y. Onel, F. Ozok, S. Sen, J. Wetzel, T. Yetkin, K. Yi

Johns Hopkins University, Baltimore, USA

B.A. Barnett, B. Blumenfeld, A. Bonato, C. Eskew, D. Fehling, G. Giurgiu, A.V. Gritsan, Z.J. Guo, G. Hu, P. Maksimovic, S. Rappoccio, M. Swartz, N.V. Tran, A. Whitbeck

The University of Kansas, Lawrence, USA

P. Baringer, A. Bean, G. Benelli, O. Grachov, M. Murray, V. Radicci, S. Sanders, J.S. Wood, V. Zhukova

Kansas State University, Manhattan, USA

D. Bandurin, T. Bolton, I. Chakaberia, A. Ivanov, K. Kaadze, Y. Maravin, S. Shrestha, I. Svintradze, Z. Wan

Lawrence Livermore National Laboratory, Livermore, USA

J. Gronberg, D. Lange, D. Wright

University of Maryland, College Park, USA

D. Baden, M. Boutemur, S.C. Eno, D. Ferencek, N.J. Hadley, R.G. Kellogg, M. Kirn, A. Mignerey, K. Rossato, P. Rumerio, F. Santanastasio, A. Skuja, J. Temple, M.B. Tonjes, S.C. Tonwar, E. Twedt

Massachusetts Institute of Technology, Cambridge, USA

B. Alver, G. Bauer, J. Bendavid, W. Busza, E. Butz, I.A. Cali, M. Chan, D. D'Enterria, P. Everaerts, G. Gomez Ceballos, M. Goncharov, K.A. Hahn, P. Harris, Y. Kim, M. Klute, Y.-J. Lee, W. Li, C. Loizides, P.D. Luckey, T. Ma, S. Nahn, C. Paus, C. Roland, G. Roland, M. Rudolph, G.S.F. Stephans, K. Sumorok, K. Sung, E.A. Wenger, B. Wyslouch, S. Xie, Y. Yilmaz, A.S. Yoon, M. Zanetti

University of Minnesota, Minneapolis, USA

P. Cole, S.I. Cooper, P. Cushman, B. Dahmes, A. De Benedetti, P.R. Duderio, G. Franzoni, J. Haupt, K. Klapoetke, Y. Kubota, J. Mans, D. Petyt, V. Rekovic, R. Rusack, M. Sasseville, A. Singovsky

University of Mississippi, University, USA

L.M. Cremaldi, R. Godang, R. Kroeger, L. Perera, R. Rahmat, D.A. Sanders, P. Sonnek, D. Summers

University of Nebraska-Lincoln, Lincoln, USA

K. Bloom, S. Bose, J. Butt, D.R. Claes, A. Dominguez, M. Eads, J. Keller, T. Kelly, I. Kravchenko, J. Lazo-Flores, C. Lundstedt, H. Malbouisson, S. Malik, G.R. Snow

State University of New York at Buffalo, Buffalo, USA

U. Baur, I. Iashvili, A. Kharchilava, A. Kumar, K. Smith, M. Strang, J. Zennamo

Northeastern University, Boston, USA

G. Alverson, E. Barberis, D. Baumgartel, O. Boeriu, S. Reucroft, J. Swain, D. Wood, J. Zhang

Northwestern University, Evanston, USA

A. Anastassov, A. Kubik, R.A. Ofierzynski, A. Pozdnyakov, M. Schmitt, S. Stoynev, M. Velasco, S. Won

University of Notre Dame, Notre Dame, USA

L. Antonelli, D. Berry, M. Hildreth, C. Jessop, D.J. Karmgard, J. Kolb, T. Kolberg, K. Lannon, S. Lynch, N. Marinelli, D.M. Morse, R. Ruchti, J. Slaunwhite, N. Valls, J. Warchol, M. Wayne, J. Ziegler

The Ohio State University, Columbus, USA

B. Bylsma, L.S. Durkin, J. Gu, P. Killewald, T.Y. Ling, G. Williams

Princeton University, Princeton, USA

N. Adam, E. Berry, P. Elmer, D. Gerbaudo, V. Halyo, A. Hunt, J. Jones, E. Laird, D. Lopes Pegna, D. Marlow, T. Medvedeva, M. Mooney, J. Olsen, P. Piroué, D. Stickland, C. Tully, J.S. Werner, A. Zuranski

University of Puerto Rico, Mayaguez, USA

J.G. Acosta, X.T. Huang, A. Lopez, H. Mendez, S. Oliveros, J.E. Ramirez Vargas, A. Zatzerklyaniy

Purdue University, West Lafayette, USA

E. Alagoz, V.E. Barnes, G. Bolla, L. Borrello, D. Bortoletto, A. Everett, A.F. Garfinkel, Z. Gecse, L. Gutay, M. Jones, O. Koybasi, A.T. Laasanen, N. Leonardo, C. Liu, V. Maroussov, P. Merkel, D.H. Miller, N. Neumeister, K. Potamianos, I. Shipsey, D. Silvers, H.D. Yoo, J. Zablocki, Y. Zheng

Purdue University Calumet, Hammond, USA

P. Jindal, N. Parashar

Rice University, Houston, USA

V. Cuplov, K.M. Ecklund, F.J.M. Geurts, J.H. Liu, J. Morales, B.P. Padley, R. Redjimi, J. Roberts

University of Rochester, Rochester, USA

B. Betchart, A. Bodek, Y.S. Chung, P. de Barbaro, R. Demina, H. Flacher, A. Garcia-Bellido, Y. Gotra, J. Han, A. Harel, D.C. Miner, D. Orbaker, G. Petrillo, D. Vishnevskiy, M. Zielinski

The Rockefeller University, New York, USA

A. Bhatti, L. Demortier, K. Goulianos, K. Hatakeyama, G. Lungu, C. Mesropian, M. Yan

Rutgers, the State University of New Jersey, Piscataway, USA

O. Atramentov, Y. Gershtein, R. Gray, E. Halkiadakis, D. Hidas, D. Hits, A. Lath, K. Rose, S. Schnetzer, S. Somalwar, R. Stone, S. Thomas

University of Tennessee, Knoxville, USA

G. Cerizza, M. Hollingsworth, S. Spanier, Z.C. Yang, A. York

Texas A&M University, College Station, USA

J. Asaadi, R. Eusebi, J. Gilmore, A. Gurrola, T. Kamon, V. Khotilovich, R. Montalvo, C.N. Nguyen, J. Pivarski, A. Safonov, S. Sengupta, D. Toback, M. Weinberger

Texas Tech University, Lubbock, USA

N. Akchurin, C. Bardak, J. Damgov, C. Jeong, K. Kovitanggoon, S.W. Lee, P. Mane, Y. Roh, A. Sill, I. Volobouev, R. Wigmans, E. Yazgan

Vanderbilt University, Nashville, USA

E. Appelt, E. Brownson, D. Engh, C. Florez, W. Gabella, W. Johns, P. Kurt, C. Maguire, A. Melo, P. Sheldon, J. Velkovska

University of Virginia, Charlottesville, USA

M.W. Arenton, M. Balazs, M. Buehler, S. Conetti, B. Cox, R. Hirosky, A. Ledovskoy, C. Neu, R. Yohay

Wayne State University, Detroit, USA

S. Gollapinni, K. Gunthoti, R. Harr, P.E. Karchin, M. Mattson, C. Milstène, A. Sakharov

University of Wisconsin, Madison, USA

M. Anderson, M. Bachtis, J.N. Bellinger, D. Carlsmith, S. Dasu, S. Dutta, J. Efron, L. Gray, K.S. Grogg, M. Grothe, R. Hall-Wilton¹, M. Herndon, P. Klabbers, J. Klukas, A. Lanaro, C. Lazaridis, J. Leonard, D. Lomidze, R. Loveless, A. Mohapatra, G. Polese, D. Reeder, A. Savin, W.H. Smith, J. Swanson, M. Weinberg

†: Deceased

- 1: Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland
- 2: Also at Universidade Federal do ABC, Santo Andre, Brazil
- 3: Also at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France
- 4: Also at Fayoum University, El-Fayoum, Egypt
- 5: Also at Soltan Institute for Nuclear Studies, Warsaw, Poland
- 6: Also at Université de Haute-Alsace, Mulhouse, France
- 7: Also at Moscow State University, Moscow, Russia
- 8: Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary
- 9: Also at Eötvös Loránd University, Budapest, Hungary
- 10: Also at University of California, San Diego, La Jolla, USA
- 11: Also at Tata Institute of Fundamental Research - HECR, Mumbai, India
- 12: Also at University of Visva-Bharati, Santiniketan, India
- 13: Also at Facoltà Ingegneria Università di Roma "La Sapienza", Roma, Italy
- 14: Also at Università della Basilicata, Potenza, Italy
- 15: Also at Laboratori Nazionali di Legnaro dell' INFN, Legnaro, Italy
- 16: Also at California Institute of Technology, Pasadena, USA
- 17: Also at Faculty of Physics of University of Belgrade, Belgrade, Serbia
- 18: Also at Université de Genève, Geneva, Switzerland
- 19: Also at Scuola Normale e Sezione dell' INFN, Pisa, Italy
- 20: Also at INFN Sezione di Roma; Università di Roma "La Sapienza", Roma, Italy
- 21: Also at University of Athens, Athens, Greece
- 22: Also at The University of Kansas, Lawrence, USA
- 23: Also at Institute for Theoretical and Experimental Physics, Moscow, Russia
- 24: Also at Paul Scherrer Institut, Villigen, Switzerland
- 25: Also at Vinca Institute of Nuclear Sciences, Belgrade, Serbia
- 26: Also at Adiyaman University, Adiyaman, Turkey

- 27: Also at Mersin University, Mersin, Turkey
- 28: Also at Izmir Institute of Technology, Izmir, Turkey
- 29: Also at Kafkas University, Kars, Turkey
- 30: Also at Suleyman Demirel University, Isparta, Turkey
- 31: Also at Ege University, Izmir, Turkey
- 32: Also at Rutherford Appleton Laboratory, Didcot, United Kingdom
- 33: Also at INFN Sezione di Perugia; Università di Perugia, Perugia, Italy
- 34: Also at KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary
- 35: Also at Institute for Nuclear Research, Moscow, Russia
- 36: Also at Istanbul Technical University, Istanbul, Turkey